

THERMAL RESPONSE TO FIRE OF UNIFORMLY INSULATED STEEL MEMBERS: BACKGROUND AND VERIFICATION OF THE FORMULATION RECOMMENDED BY CHINESE CODE CECS200

Guo-qiang Li^{1,2} and Chao Zhang^{2,*}

¹Sate Key Laboratory for Disaster Reduction in Civil Engineering, 1239 Siping Road, Shanghai, China

²College of Civil Engineering, Tongji University, 1239 Siping Road, Shanghai 200092, China

*(Corresponding author: E-mail: 08_chao_zhang@tongji.edu.cn)

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ABSTRACT: Currently, Chinese technical code for fire safety of steel structures in buildings (CECS200) is being revised. This paper intends to give the background of the recommended formulation in CECS200 for temperature calculation of uniformly insulated steel members. Analytical formulations recommended by other codes including EC3, ECCS, SFPE handbook, etc., and FE method are used for comparison and verification. Theoretical derivation of the exact analytical formulations is also given. Two boundary conditions at the fire-insulation interface have been discussed in the derivation. By investigating the steel temperature of three insulated sections with different section factors which are protected by typical fire protection materials, the formulation recommended by CECS200 is proved to be reasonable and simple for engineering usage.

Keywords: Fire resistance, uniformly insulated steel members, 1D model, temperature calculation, analytical formulations, FEM

1. INTRODUCTION

Bare Steel is frangible to fire by its high conductivity and low specific heat capacity. As a result, insulation is always required for steel members to ensure the stability of the building in fire condition. Historically, the most common methods of insulation have included concrete encasement, envelopment in gypsum plaster or gypsum board and spray application of light weight cementitious or mineral fiber spray-applied fire resistive materials (SFRMs). In recent years, intumescent coatings have been widely used in fire protection engineering. Unlike the traditional fire protection materials like concrete, gypsum, SFRMs, etc., the intumescent coatings are reactive materials that they are 'inert' at low temperatures but swell to provide a charred layer of low conductivity materials at temperatures of approximately 200-250°C[1].

Traditionally, the thickness of insulation needed on steel members is determined by standard fire tests [2-3]. In standard fire tests, after a specified period of time, if the insulated member withstands the fire exposure without exceeding any of the endpoint failure criteria, the insulation is sufficient and the thickness of the insulation is considered to be the limit value of the protection. The standard fire conditions in the tests are represented in terms of the standard gas temperature history specified in different standards like ASTM E119 [2], ISO834 [3], etc... The widely used endpoint failure criteria in fire resistance test standard is that the maximum mean steel temperature must be lower than the critical temperature which is the temperature that causes structure collapse in a fire situation.

As an alternative to standard fire testing, analytical methods have been developed to evaluate the fire protection materials [4-7]. Analytical methods require solution of 2D (ignoring temperature gradient along the length of the member) transient heat transfer diffusion equation [8], which is usually complicated. In practice, a simplified 1D condensed heat transfer model, as shown in Figure 1, based on lumped capacitance concept has been adopted by different codes [4-7]. This

concept assumes that the temperature distribution is uniform inside the entire steel section. Using this 1D model, various mathematical techniques including separation of variables [9], Laplace transform [10] and Green's function approach [11] have been developed to give analytical formulations for uniformly insulated steel members in fire. There are two boundary conditions, which are Dirichlet and Neumann boundary conditions, at the fire-insulation interface for the solution of the 1D diffusion equation. In reference [9], Dirichlet boundary condition has been assumed and the derived formulation has been adopted by Eurocode 3[4]. In references [10] and [11], both Dirichlet and Neumann boundary conditions can be considered. Assuming Dirichlet boundary condition, Silva [12] derives a simple formulation of the problem which has been recommended for the revision of the Brazilian Standard 14323 [13].

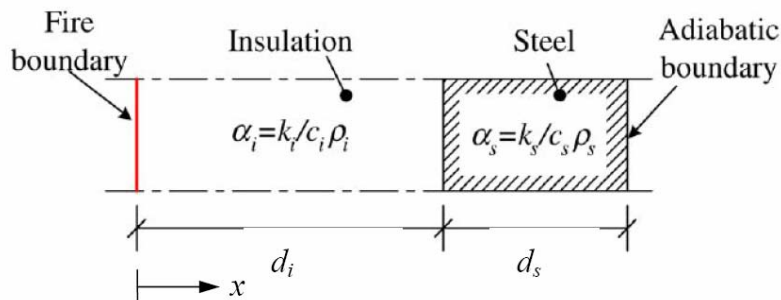


Figure 1. 1D Heat Transfer Model

Currently, Chinese technical code for fire safety of steel structure in buildings (CECS200) [7] is being revised. This paper intends to give the background of the formulation given by CECS200 [7] for temperature calculation of uniformly insulated steel members. The exact analytical derivations of the problem are also presented. Three steel I sections with different section factors insulated by typical nonreactive fire protection materials which are normal weight concrete (NWC), gypsum board and SFRMs are investigated. The mean steel temperature calculated by CECS200 [7], EC3 [4], ECCS [5], SFPE handbook [6] and Silva [12] are compared with the numerical results from the popular FEM program ANSYS [14]. ANSYS is power to solve transient, non-linear heat transfer problems. Its capacity of application in fire condition has been verified by experiments [15]. The same thermal elements used in reference [15] are adopted in this paper. In Appendix, the FEM model is further verified. Detail information about the FEM model is also given in the Appendix.

2. THEORETICAL BACKGROUND

2.1 Governing Equation and Boundary Conditions

The governing heat transfer equation for the 1D model in Figure 1 is given by

$$\alpha_i \frac{\partial^2 T(x,t)}{\partial x^2} - \frac{\partial T(x,t)}{\partial t} = 0 \quad (1)$$

where $\alpha_i = k_i / \rho_i c_i$ is the thermal diffusivity; k_i is the thermal conductivity; and $\rho_i c_i$ is the thermal capacity, in which ρ_i , c_i are the density and specific heat of the insulation material, respectively.

By lumped capacitance concept, the boundary condition at the steel-insulation interface is given by

$$-k_i \frac{\partial T(d_i, t)}{\partial x} = \frac{c_s \rho_s}{A_p / V} \frac{\partial T(d_i, t)}{\partial t} \quad (2)$$

$$T_s(t) = T(d_i, t) \quad (3)$$

At the fire-insulation interface, the Neumann boundary condition is

$$-k_i \frac{\partial T(0, t)}{\partial x} = (h_c + h_r)[T_g(t) - T(0, t)] \quad (4)$$

where, $h_c = 25 \text{ W}/(\text{m}^2\text{K})$ is the convective heat transfer coefficient; h_r is the radiative heat transfer coefficient, given by

$$h_r = \sigma \varepsilon_{res} [(T_g(t) + 273)^2 + (T(0, t) + 273)^2] \times [T_g(t) + 273 + T(0, t) + 273] \quad (5)$$

where, $\sigma = 5.67 \times 10^{-8} \text{ W}/(\text{m}^2\text{K}^4)$ is the Stefan-Boltzmann constant; and ε_{res} is the resultant emissivity for radiation heat transfer, given by

$$\varepsilon_{res} = \frac{1}{1/\varepsilon_g + 1/\varepsilon_m - 1} \approx \varepsilon_g \varepsilon_m \quad (6)$$

where ε_g and ε_m are the emissivity of the fire and insulation surface respectively. The emissivity of a real surface is a function of the surface temperature [8, 16-17], but in present codes [4-7], constant values of emissivity are adopted. In this paper, constant values with $\varepsilon_g=1.0$ and $\varepsilon_m=0.7$ recommended in EC4 [18] are used.

The Neumann boundary condition at the fire-insulation interface is complex. In practice, Dirichlet boundary condition has been widely adopted in deriving temperature formulations for uniformly insulated steel members [4, 7, 9, 12], which is given by

$$T(0, t) = T_g(t) \quad (7)$$

Eq. 7 assumes that there is no heat loss through surface convection and radiation, thus the insulation can effectively ‘block’ the thermal energy at the fire-insulation interface [19]. This assumption gives good prediction of steel temperature for insulation with low density and low conductivity but yields conservative results for insulation with high density and high conductivity [16]. In reference [16], two insulation materials, ceramic fibre blanket and concrete, have been studied, and temperature-dependent emissivity has been used. Using Green’s function, Wang et al. [11] shows that for steel sections insulated by gypsum and concrete, the Dirichlet boundary condition will yield higher steel temperature than the real ones using Neumann boundary condition.

2.2 Exact Analytical Derivations

2.2.1 Separation of Variables

Introducing a state variable defined by

$$\theta(x,t) = T(x,t) - T_0 \quad (8)$$

Then, the governing equation Eq. 1 and the boundary conditions Eqs. 2 and 7 are expressed as

$$\alpha_i \frac{\partial^2 \theta(x,t)}{\partial x^2} - \frac{\partial \theta(x,t)}{\partial t} = 0 \quad (9)$$

$$-k_i \frac{\partial \theta(d_i,t)}{\partial x} = \frac{c_s \rho_s}{A_p / V} \frac{\partial \theta(d_i,t)}{\partial t} \quad (10)$$

$$\theta(0,t) = 0 \quad (11)$$

Solving the correlated Eqs. 9 to 12 with the following initial condition

$$\theta(x,0) = \theta_0, \quad t = 0 \quad (12)$$

we get

$$\theta(x,t) / \theta_0 = \sum_{n=1}^{\infty} C_n \exp\left(-\frac{\alpha_i \xi_n^2}{d_i^2} t\right) \sin\left(\xi_n \frac{x}{d_i}\right) \quad (13)$$

From Eq. 3, we get $\theta_s(t) = \theta(d_i,t)$, then

$$\theta_s(t) / \theta_0 = \sum_{n=1}^{\infty} C_n \exp\left(-\frac{\alpha_i \xi_n^2}{d_i^2} t\right) \sin(\xi_n) \quad (14)$$

In which, ξ_n is obtained by solving the following transcendental equation

$$\xi_n \tan \xi_n = \mu \quad (15)$$

where,

$$\mu = \frac{c_i \rho_i}{c_s \rho_s} d_i (A_p / V) \quad (16)$$

The coefficients C_n in Eq. 14 are obtained by

$$C_n = \frac{2(\xi_n^2 + \mu^2)}{\xi_n(\xi_n^2 + \mu^2 + \mu)} \quad (17)$$

In the case that μ is small, Eq. 14 can be approximated as

$$\theta_s(t)/\theta_0 = \begin{cases} 0 & t \leq t_d \\ \exp(-\frac{t-t_d}{\tau}) & t > t_d \end{cases} \quad (18)$$

where,

$$t_d = \mu\tau/8, \quad \tau = \frac{c_s \rho_s d_i}{(A_p/V)k_i} (1 + \frac{\mu}{3}) \quad (19)$$

Assuming the thermal properties are constant, temperature response can be calculated for the insulated steel member subject to a time-varying boundary condition, using the principle of superposition (Duhamel's theorem), as

$$T_s(t) = \int_0^t T(t-\zeta) d(1-\theta_s(\zeta)/\theta_0) \quad (20)$$

where T is the time-varying heating curve. The ISO834 standard fire curve can be approximated by a sum of exponential terms as

$$T = \sum_{j=0}^3 B_j \exp(-\beta_j t) \quad (21)$$

where B_j and β_j are as given in Table 1.

Table 1. B_j, β_j for ISO834 Standard Curve

j	0	1	2	3
$B(^{\circ}\text{C})$	1325	-430	-270	-625
$\beta(\text{h}^{-1})$	0	0.2	1.7	19

Substituting Eqs. 14 and 21 into Eq. 20 we get the steel temperature

$$T_s = \sum_{n=1}^{\infty} \sum_{j=0}^3 \frac{B_j C_n \sin(\xi_n)}{1 - \beta_j d_i^2 / (\alpha_i \xi_n^2)} \times [\exp(-\beta_j t) - \exp(-\frac{\alpha_i \xi_n^2}{d_i^2} t)] \quad (22)$$

Substituting Eqs. 18 and 21 into Eq. 20, we get the approximate steel temperature, for $t > t_d$,

$$T_s = \sum_{j=0}^3 \frac{B_j}{1 - \beta_j \tau} \times \left[\exp(-\beta_j (t-t_d)) - \exp(-\frac{t-t_d}{\tau}) \right] \quad (23)$$

To consider time-varying material properties, an approximate temperature time derivative is derived by Wickström[9] from Eq. 23 as

$$\frac{dT_s}{dt} = \frac{T - T_s}{\tau} - (e^{\mu/10} - 1) \frac{dT}{dt} \quad (24)$$

Eq. 24 is adopted by EC3 [4], and the increment of steel temperature ΔT_s within a time interval Δt is given by

$$\Delta T_s = \frac{k_i A_p / V}{c_s \rho_s d_i} \frac{T_g - T_s}{(1 + \mu/3)} \Delta t - (e^{\mu/10} - 1) \Delta T_g \quad (\Delta T_s \geq 0, \text{ if } \Delta T_g > 0) \quad (25)$$

where, $\Delta t \leq 30\text{s}$ for insulated steel members.

2.2.2 Laplace Transform

An alternative solution of the problem using Laplace transfer is given by ECCS [5] as

$$\frac{dT_s}{dt} = A'(T_g - T_s) - B' \frac{dT_g}{dt} \quad (26)$$

where

$$A' = \frac{1}{\left(\frac{c_s \rho_s}{A_p / V}\right) \left(\frac{d_i}{k_i} + \frac{1}{h_c + h_r}\right) \left(1 + \frac{\mu}{N}\right)} \quad (27)$$

$$B' = b / (1 + N / \mu) \quad (28)$$

with N and b as weighting factors. Certainly, for limiting case $(h_c + h_r) \rightarrow \infty$ Neumann boundary is equivalent to Dirichlet boundary. At this case,

$$b = \frac{1 + \mu/4}{2(1 + 5\mu/8)} \quad (29)$$

$$N = 2(b + 1) \quad (30)$$

2.3 Simple Derivation in CECS200

In Chinese Code CECS200 [7], the following assumptions are made in deriving the formulation for temperature calculation of insulated steel members,

- (1) Dirichlet boundary is safely assumed at fire-insulation interface;
- (2) The temperature distribution within the insulation is linear; and
- (3) The temperature distribution within the steel is uniform.

At time increment Δt , the total energy transferred to the steel is

$$\Delta Q = \frac{k_i}{d_i} [T(t) - T_s(t)] A_p \Delta t \quad (31)$$

The energy absorbed by the steel is

$$\Delta Q_s = c_s \rho_s V [T_s(t + \Delta t) - T_s(t)] \quad (32)$$

The energy absorbed by the insulation is

$$\Delta Q_i = \frac{T_s(t + \Delta t) - T_s(t) + \Delta T_g}{2} c_i \rho_i A_p d_i \quad (33)$$

By energy balance, we have

$$\Delta Q = \Delta Q_s + \Delta Q_i \quad (34)$$

Substituting Eqs. 31, 32 and 33 into Eq. 34 and ignoring the secondary term, we get the formulation given by CECS200 [7] as

$$\Delta T_s = \frac{k_i A_p / V}{c_s \rho_s d_i} \frac{T_g - T_s}{(1 + \mu / 2)} \Delta t \quad (35)$$

3. COMPARISON OF DIFFERENT FORMULATIONS

3.1 Different Formulations

Eq. 25, Eq. 26 and Eq. 35 give the formulations recommended by EC3 [4], ECCS [5] and CECS200 [7] respectively. In SFPE handbook [6] different formulations are recommended based on the value of μ . For $\mu > 1/2$, the same formulation in CECS200 or Eq. 35 is recommended; and for $\mu \leq 1/2$,

$$\Delta T_s = \frac{k_i A_p / V}{c_s \rho_s d_i} (T_g - T_s) \Delta t \quad (36)$$

The formulation recommended by Silva [13] is

$$\Delta T_s = \frac{k_i A_p / V}{c_s \rho_s d_i} \frac{T_g - T_s}{(1 + \mu / 4)} \Delta t - \frac{\Delta T_g}{4 / \mu + 1} \quad (37)$$

The formulation recommend by Pettersson et al. [20] is

$$\Delta T_s = \frac{A_p / V}{c_s \rho_s} \frac{1}{\left(\frac{d_i}{k_i} + \frac{1}{h_c + h_r}\right) \left(1 + \frac{\mu}{2}\right)} \frac{T_g - T_s}{\mu} \Delta t - \frac{\Delta T_g}{\frac{2}{\mu} + 1} \quad (38)$$

3.2 Material and Geometric Properties

Figure 2 shows the schematics of the analyzed section. Totally 3 typical steel I section with different section factors are investigated. The geometric properties of the sections are given in Table 2. The thickness of all sections is 50mm.

Table 3 gives the thermal properties of fire protection materials [21] which are NWC, gypsum board and SFRMs. Temperature-independent properties are investigated. The thermal properties of the steel are $\rho_s=7850 \text{ kg/m}^3$ and $c_s=600 \text{ J/(kg}^\circ\text{C)}$.

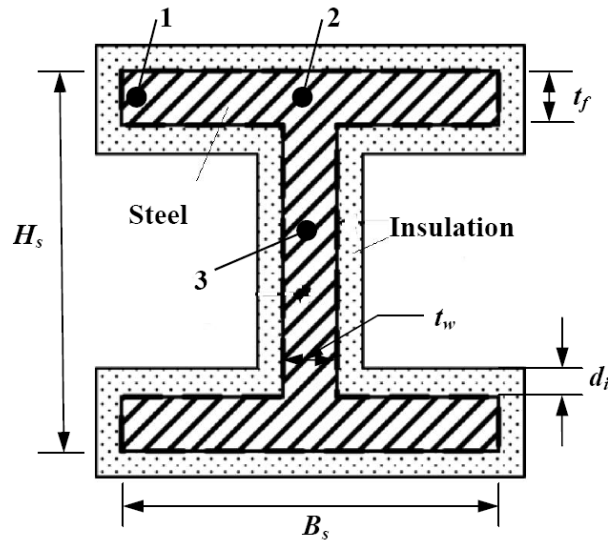


Figure 2. Schematics of Analyzed Section

Table 2. Geometric Properties of Analyzed Sections

Section name	B_s (mm)	H_s (mm)	t_w (mm)	t_f (mm)	A_p/V (m^{-1})	d_i (mm)
#1	200	200	8	12	190.72	50
#2	300	300	10	15	152.14	50
#3	400	400	15	25	93.86	50

Table 3. Thermal Properties of Insulation Materials

	ρ_i (kg/m^3)	c_i J/($kg \text{ } ^\circ C$)	k_i W/($m \text{ } ^\circ C$)
SFRM	250	800	0.12
gypsum board	800	1700	0.2
NWC	2300	1200	1.6

3.3 Results and Discussions

Figure 3 to Figure 5 show the steel temperature-time curves of insulated sections resulting from different methods. The steel temperatures obtained with FEM are taken as the average value of the temperatures measured at the location of points 1, 2 and 3, as shown in Figure 2.

When calculating by Eqs. 25, 26, 37 and 38, the increments of the steel temperature are negative at early heating process, which is illegal to physics law. To avoid this, in practice we assume the steel temperature increments as zero if the values calculated by those equations are negative.

When calculating by Eqs. 26 and 38, the values of h_r are required. From Eq. 5 we know in order to get h_r , the value of the surface temperature of the insulation, $T(0, t)$ should be known beforehand. However, $T(0, t)$ is an unknowable variable. In practice when calculating h_r we safely take fire temperature as the surface temperature that $T(0, t) \approx T_g(t)$. This treatment has been used by Wong and Ghojel [16] in their analysis.

As shown in Figure 3, when insulated by SFRMs, for all sections steel temperature – time curves given by EC3 [4], ECCS [5], CECS200 [7] and Silva [12] fit very well. Curves given by SFPE handbook [6] and Pettersson et al. [20] deviate from curves given by other four formulations. Comparatively, SFPE handbook [6] over-predicts the steel temperature and Pettersson et al. [20] under-predicts the steel temperature.

As shown in Figure 4, when insulated by gypsum board, for all sections temperature – time curves given by EC3 [4], ECCS [5] and Silva [12] fit well. CECS200 [7] and SFPE handbook [6] give the same results which fit well with the results given by EC3, ECCS and Silva at late heating stage but safely deviate from those results at early and mid heating stage. Comparing with other formulations, Pettersson et al. [20] under-predicts the steel temperature during the whole heating period.

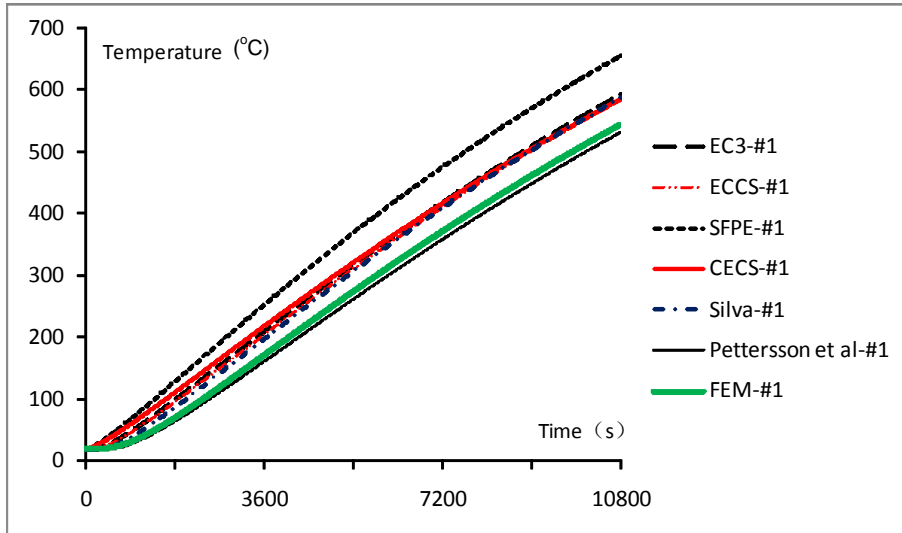
As shown in Figure 5, when insulated by NWC, for all sections temperature – time curves given by EC3 [4], ECCS [5] and CECS200 [7] and SFPE handbook [6] fit very well. Silva [12] gives similar but slightly conservative results. Pettersson et al. [20] gives unsafe results.

As shown in Figures. 3-5, comparing with the ‘exact’ FEM results, except Pettersson et al. [20] under-predicts the steel temperature all other analytical methods will over-predict the steel temperature. Results also show section factors have little effect on the relationship between different curves.

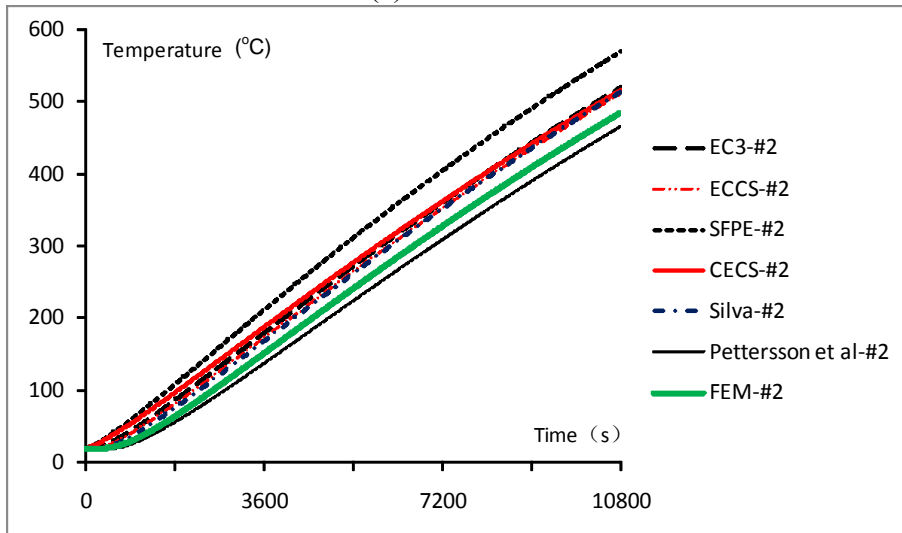
4. CONCLUSIONS

Based on the results of this study, the following conclusions can be drawn:

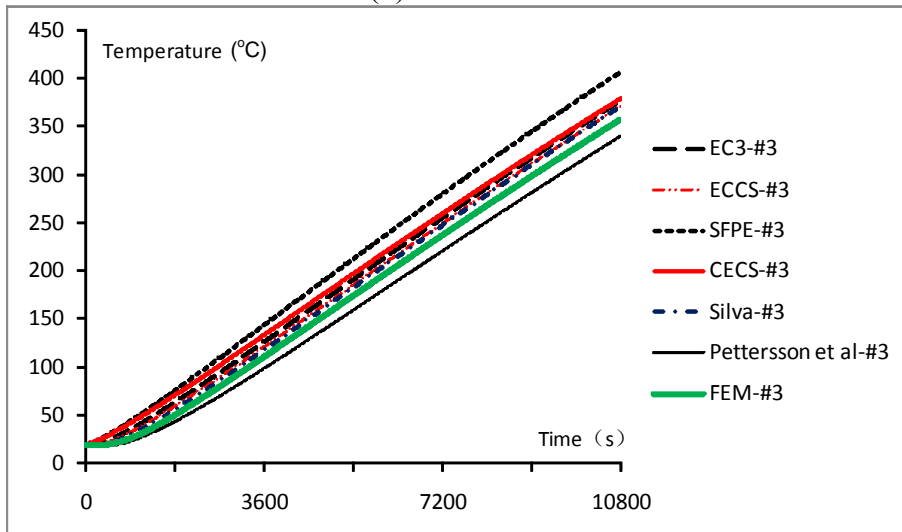
- Comparing with the ‘exact’ FEM method, except formulation recommended by Pettersson et al. [20] other formulations will yield conservative results. This is because in the derivation of those formulations, Dirichlet boundary has been assumed at the fire-insulation interface.
- Formulations recommended by EC3 [4] and ECCS [5] nearly give the same results. Results given by formulations recommended by CECS200 [7] fit well with the results given by formulations recommended by EC3 [4], ECCS [5] and Silva [12].
- When calculating by formulation recommended by CECS200 [7], there will not occur unreasonable situation that steel temperature increment is negative in early heating state. Considering its simple form, formulation recommended by CECS200 [7] is preferable for engineering usage.



(a) 1# section

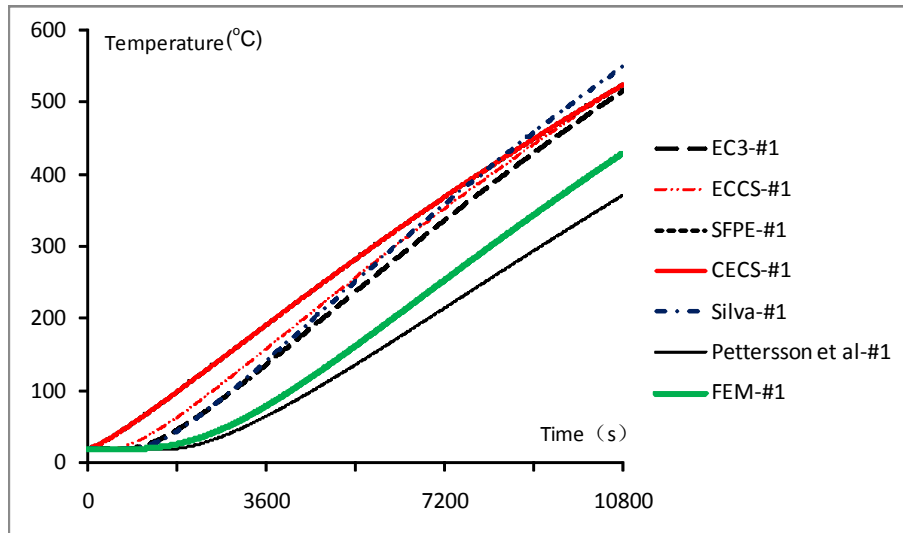


(b) 2# section

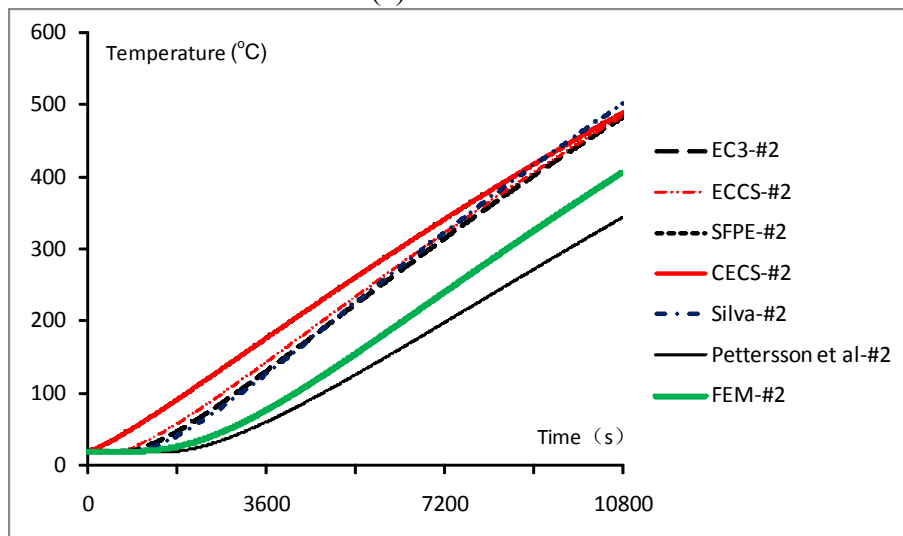


(c) 3# section

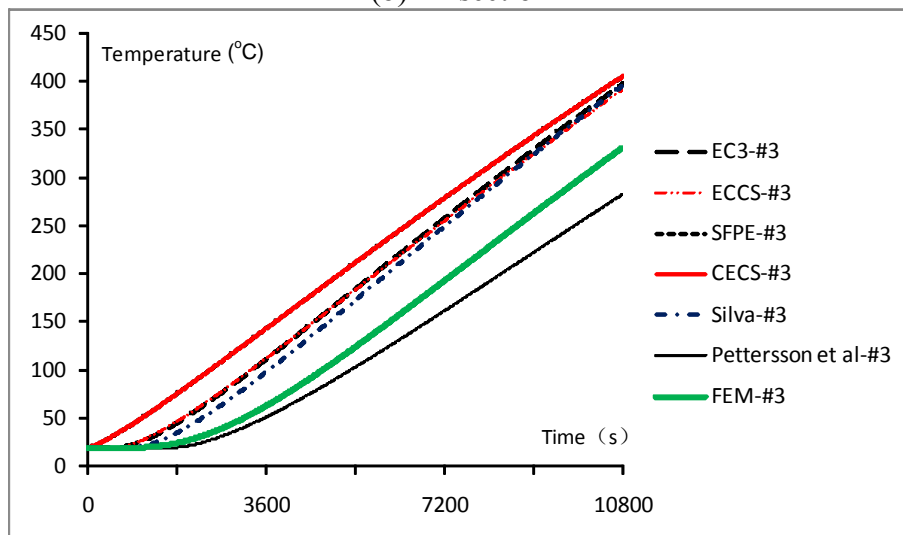
Figure 3. Results of Steel Temperature when Insulated by SFRMs



(a) 1# section

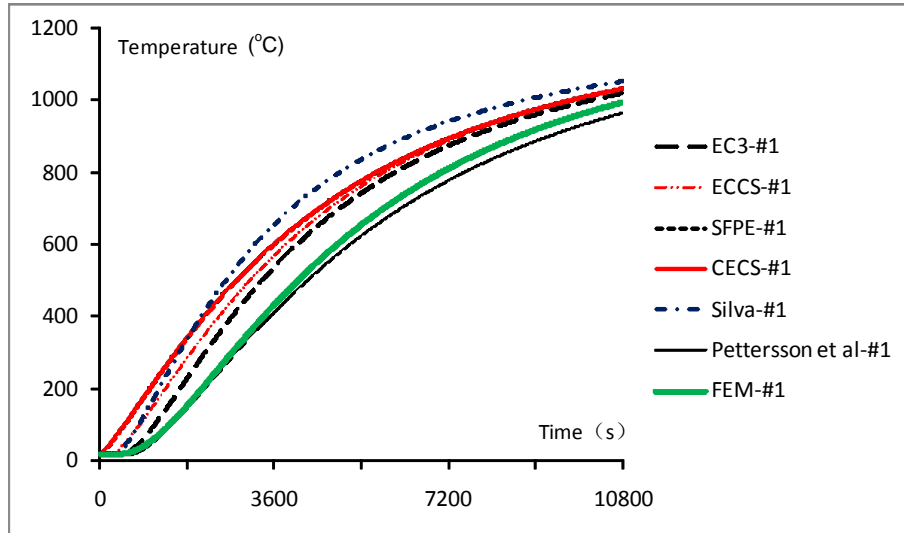


(b) 2# section

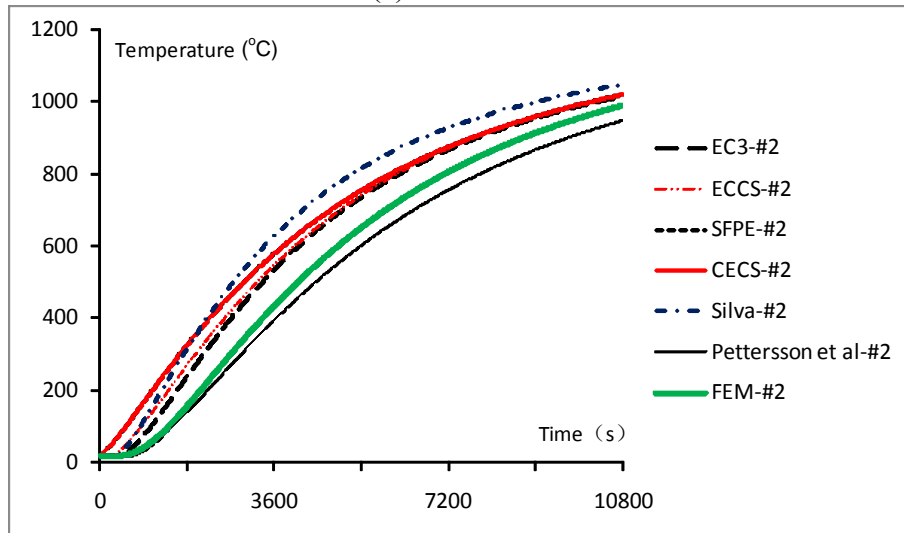


(c) 3# section

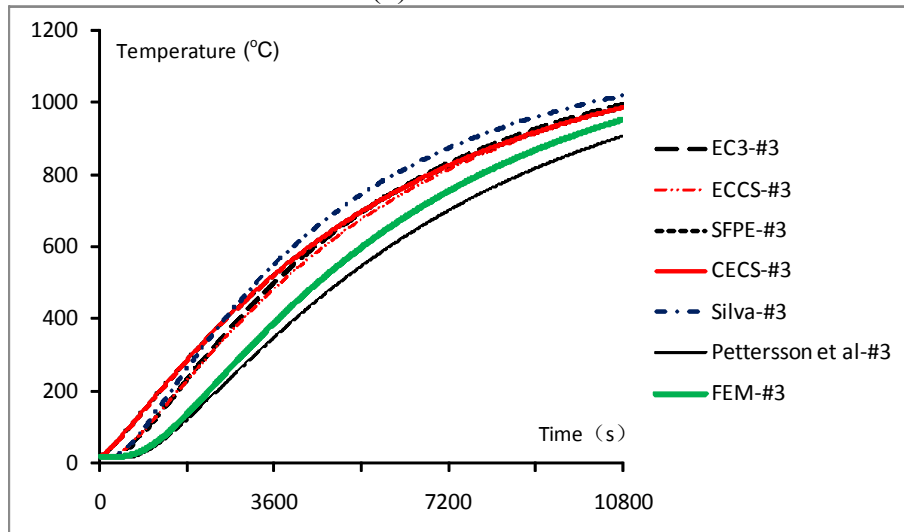
Figure 4. Results of Steel Temperature when Insulated by Gypsum Board



(a) 1# section



(b) 2# section



(c) 3# section

Figure 5. Results of Steel Temperature when Insulated by NWC

APPENDIX

In numerical simulation by ANSYS, 2D thermal solid element PLANE55 and thermal surface effect element SURF151 are adopted.

PLANE55 can be used as a plane element or as an axisymmetric ring element with a two-dimensional thermal conduction capacity. The element has four nodes with a single degree of freedom, temperature, at each node. The element is applicable to a two-dimensional, steady-state or transient thermal analysis. It can also compensate for mass transport heat flow from a constant velocity field.

SURF151 may be used for various load and surface effect applications. It may be overlaid onto a face of any 2D thermal solid element (except axisymmetric harmonic elements). The element is applicable to two-dimensional thermal analysis. Various loads and surface effects may exist simultaneously. For example, SURF151 can be overlaid onto the surface of PLANE55 to simulate the effect of thermal radiation from ambient air to steel section.

Problem: A ceramic wall is initially uniform in temperature at 20 °C and has a thickness of 30mm. It is suddenly exposed to a radiation source on the right side at 1000°C. The left side is exposed to room air at 20 °C with a radiation surrounding temperature of 20 °C. Properties of the ceramic are $k = 3.0 \text{ W/(m } ^\circ\text{C)}$, $\rho = 1600 \text{ kg/m}^3$, and $c = 800 \text{ J/(kg } ^\circ\text{C)}$. Radiation heat transfer with the surroundings at T_r (in °C) may be calculated from

$$q_r = 0.8\sigma A[(T + 273)^2 - (T_r + 273)^2]$$

The convection heat transfer coefficient from the left side of the plate is given by

$$h = 1.92\Delta T^{1/4}$$

Determine the temperature distribution in the plate after 15,30,45,60,90,120, and 150 s.

Holman[8] solves the above problem by finite difference method (FDM). Here we solve the same problem by FEM using ANSYS. Figure. 6 shows the results from FDM and FEM. It finds the results from FEM fit well with the results given by Holman using FDM, which verify the validation of the FEM in this paper.

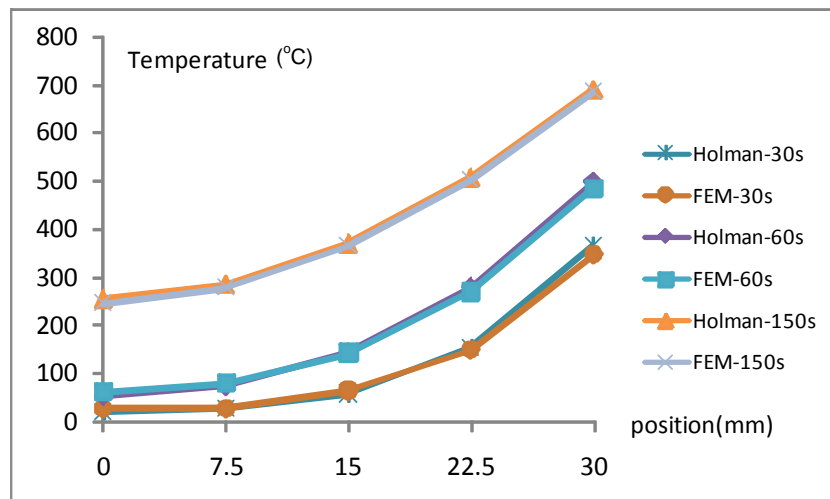


Figure 6. Comparison of Results by FEM and Holman [8]

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